Extended Positron-Trapping Defects in the Eu³⁺-Doped BaGa₂O₄ Ceramics Studied by Positron Annihilation Lifetime Method

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The BaGa₂O₄ ceramics doped with Eu³⁺ ions (1, 3 and 4 mol%) are obtained by solid-phase sintering. The evolution of defect-related extended free volumes in the BaGa₂O₄ ceramics due to the increase of the content of Eu³⁺ ions has been studied using the positron annihilation lifetime (PAL) spectroscopy technique. It is established that the increase in the number of Eu³⁺ ions in the basic BaGa₂O₄ matrix leads to the agglomeration of free-volume defects with their subsequent fragmentation. The presence of Eu³⁺ ions results in the expansion of nanosized pores and an increase in their number with their future fragmentation.

1. Introduction

The BaGa₂O₄ ceramics are considered as a promising material for use as an insulator in optoelectronic devices,^[1,2] as a secondary coating for plasma panels,^[3–5] etc.^[6] The doping of impurities in the form of rare-earth ions leads to the expansion of the functional properties of such ceramics.^[7]

Recent studies of materials similar to $BaGa_2O_4$ ceramics in both undoped and doped forms are mainly focused on the study of their luminescent properties.^[1,8–14] These ceramics are a double oxide belonging to the quadrilateral frame topologies that exist in multiple polymorphs, which is important for luminescence studies because it exhibits luminescence without the inclusion of rare-earth ions. However, the most interesting is rare-earth-doped $BaGa_2O_4$ ceramics. Most investigations of such material are limited to X-ray diffraction (XRD), optical studies, etc.^[15–24] Nevertheless, doping of such ceramics not only leads to modification of their luminescent properties and changes in phase composition, but also causes

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structural transformation.^[25–27] Such significant transformation of the inner structure in ceramics facilitates the formation of additional defect-related free volumes.^[28–33]

Among the less conventional but promising methods to study the defect-related free volumes in solids of different structural types is positron annihilation lifetime (PAL) spectroscopy.^[34,35] This technique has long been serving as a tool for structural analysis of such functional materials as ceramics,^[36,37] glass,^[38] polymers,^[39]

nanocomposites,^[40] etc. Researchers in the field have presented a number of approaches to the analysis of PAL spectra of ceramic materials and to the proper decomposition of such spectra into different numbers of components.^[41] We have also contributed to the topic using the PAL method to study changes in free volume in spinel ceramics,^[42–44] chalcogenide glasses^[45,46] under the influence of technological modification, etc. It was shown that functional ceramics are characterized by the decomposition of the PAL spectrum into two, three, and four components depending on the developed porous structure of materials.

It was concluded that for functional ceramic materials a pair of PAL channels is enabled: positron trapping (components with shorter and medium lifetimes) and o-Ps decaying (component with longer lifetime). Generally, these processes are independent ones. However, if trapping sites emerge in the vicinity of grain boundaries neighboring free-volume pores, they can become mutually interconnected thus bringing a significant complication to the measured PAL spectra. According to the two-state positron trapping model, in oxide ceramics, the first component reflects mainly microstructural specifics of spinel ceramics with characteristic octahedral and tetrahedral vacant cation sites along with a contribution from the annihilation of p-Ps atoms which is not considered in the further analysis.^[42,47] Medium lifetime is related to the size of free-volume defects near grain boundaries and intensity reflects their amount. The third, longer component originates from the annihilation of o-Ps atoms in intrinsic nanopores of ceramics. However, the dedicated investigation of BaGa2O4 ceramics by the PAL method has not yet been conducted. Therefore, the goal of this work was to study the transformations of inner free volumes (extended defects at grain boundaries and nanopores) in BaGa2O4 ceramics doped with different amounts of Eu³⁺ ions using PAL method.





2. Results and Discussion

The best fit parameters of PAL spectra for the studied ceramics calculated within the three-components procedure are listed in **Table 1** and graphically presented in **Figure 1**. It is established that with the rising of Eu³⁺ content in the BaGa₂O₄ matrix, an increase in the lifetime τ_1 and intensity I_1 of the first short-term component is observed. It is very likely that the structure of the principal ceramics phase is improved. However, when the Eu³⁺ ions concentration increases to 4 mol.%, the opposite trend starts to prevail. The lifetime τ_2 of the second component increases and

intensity I_2 decreases with Eu³⁺ content (to 3 mol.%), while a further rise of Eu³⁺ ions (to 4 mol.%) leads to a decrease τ_2 and increase I_2 .

As noted in ref. [42,43,47], the lifetime τ_2 of the second component should be associated with the capture of positrons by defect-related free volumes. According to XRD data,^[48] undoped BaGa₂O₄ ceramics contain a large number of additional phases. As shown by the SEM method,^[48] these phases are unevenly distributed in the volume of the ceramic and mainly localized near the grain boundaries. The extracted phases play the role of special centers of positron capture in the volume of ceramics.

Table 1. Lifetimes and intensities for undoped and Eu³⁺-doped BaGa₂O₃ ceramics obtained at a three-components fitting procedure.

Sample	τ_1 [ns]	I ₁ [%]	τ ₂ [ns]	I ₂ [%]	τ_3 [ns]	I3 [%]
BaGa ₂ O ₄	$\textbf{0.200}\pm\textbf{0.002}$	83.3 ± 0.1	$\textbf{0.424} \pm \textbf{0.001}$	14.9 ± 0.1	$\textbf{2.196} \pm \textbf{0.001}$	$\textbf{1.8}\pm\textbf{0.1}$
$BaGa_2O_4 + 1 mol\% Eu^{3+}$	$\textbf{0.206} \pm \textbf{0.002}$	$\textbf{85.0}\pm\textbf{0.1}$	$\textbf{0.450} \pm \textbf{0.001}$	13.2 ± 0.1	$\textbf{2.289} \pm \textbf{0.001}$	$\textbf{1.8}\pm\textbf{0.1}$
$BaGa_2O_4 + 3 mol\% Eu^{3+}$	$\textbf{0.212} \pm \textbf{0.002}$	89.9 ± 0.1	$\textbf{0.550} \pm \textbf{0.001}$	$\textbf{7.9}\pm\textbf{0.1}$	$\textbf{2.390} \pm \textbf{0.001}$	$\textbf{2.2}\pm\textbf{0.1}$
$BaGa_2O_4 + 4 mol\% Eu^{3+}$	$\textbf{0.201} \pm \textbf{0.002}$	83.3 ± 0.1	$\textbf{0.411} \pm \textbf{0.001}$	14.4 ± 0.1	$\textbf{2.157} \pm \textbf{0.001}$	$\textbf{2.4}\pm\textbf{0.1}$



Figure 1. Dependences of lifetimes and intensities of positron annihilation lifetime (PAL) components on Eu^{3+} amount for undoped and Eu^{3+} -doped $BaGa_2O_4$ ceramics.

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Table 2. Positron trapping parameters and nanopore radius for undoped and Eu^{3+} -doped $BaGa_2O_4$ ceramics.

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Sample	$\tau_{\rm av.} \ [{\rm ns}]$	$\tau_{\rm b}~[{\rm ns}]$	$\kappa_{\rm d} \ [{\rm ns}^{-1}]$	τ ₂ _b [ns]	τ_2/τ_b	R ₃ [nm]
BaGa ₂ O ₄	0.234	0.218	0.40	0.21	1.95	0.306
$BaGa_2O_4 + 1 \text{ mol}\% \text{ Eu}^{3+}$	0.239	0.222	0.35	0.23	2.02	0.314
$BaGa_2O_4+3\ mol\%\ Eu^{3+}$	0.240	0.223	0.23	0.33	2.46	0.322
$BaGa_2O_4 + 4 mol\% Eu^{3+}$	0.232	0.218	0.37	0.19	1.89	0.302



Figure 2. Dependences of positron trapping rate in defects κ_d on Eu³⁺ amount for undoped and Eu³⁺-doped BaGa₂O₄ ceramics.

Since undoped BaGa₂O₄ ceramics contain two additional phases, positrons in such samples have statistically a better chance of being captured.

The lifetime τ_2 correlates with the size of free-volume defects, where positrons are captured and the intensity relates to their number. Thus, when the Eu³⁺ content increases to 3 mol%, agglomeration of free-volume defects occurs, while the supersaturation of the BaGa₂O₄ matrix with Eu³⁺ (up to 4 mol.%) leads to their fragmentation. The identified trends (**Table 2**) correlate with the positron trapping parameters $\tau_{av.}$, τ_b , κ_d calculated within the two-state positron trapping model.^[47] The changes are best reflected in the positron trapping rate in defects κ_d (see



Figure 4. Dependencies of lifetime τ_3 and nanopores radius R_3 calculated according to the Tao–Eldrup model for undoped and Eu³⁺-doped BaGa₂O₄ ceramics.

Figure 2). This value decreases with increasing Eu³⁺ content in ceramics and increases with its supersaturation of Eu³⁺ (up to 4 mol.%). The obtained values of τ_2/τ_b and $(\tau_2-\tau_b)$ parameters speak in favor of the similar nature of trapping sites. The characteristic size of these extended positron traps is close to that of single-double atomic vacancies.

Representatively, the evolution of defect-related free volumes near grain boundaries with increasing of Eu^{3+} ions in the $BaGa_2O_4$ ceramics is shown in **Figure 3**.

The third long-term component of the PAL spectrum (the second channel of positron annihilation with lifetime τ_3 and intensity I_3) is associated with the o-Ps decaying in nanopores.^[42] As can be seen from Table 1, lifetime τ_3 and intensity I_3 increase with Eu³⁺ content in the BaGa₂O₄ ceramics (to 3 mol%). However, when the amount of Eu³⁺ ions further rises to 4 mol%, the lifetime τ_3 decreases and the intensity I_3 continues to go up. It is obvious that the presence of Eu³⁺ ions in the ceramic matrix leads to the expansion of nanosized pores and an increase in their number. However, supersaturation of ceramics with doped ions (to 4 mol.%) results in fragmentation of nanopores. The obtained trends are reflected in the radius of nanopores R_3 calculated by the Tao-Eldrup model (**Figure 4**).

Figure 5 explains this process in a schematic fashion, outlining the main stages of the evolution.



- area change during fragmentation/agglomeration







Figure 5. Diagram explaining the evolution of nanopores in $BaGa_2O_4$ ceramics caused by Eu^{3+} doping.

Summarizing, the PAL method can also be used to assess the size of nanopores in the undoped and doped ceramic materials.

3. Conclusion

The structural features and evolution of free-volume defects in the $BaGa_2O_4$ ceramics obtained by solid-phase synthesis from the initial $BaCO_3$ and Ga_2O_3 components with the addition of different amounts of Eu_2O_3 content (1, 3, and 4 mol%) were investigated.

Additional phases in ceramics are mainly localized near the grain boundaries and create defective centers for positron capture, as studied by PAL spectroscopy. Through the analysis of the second component of PAL spectra for the undoped and ${\rm Eu}^{3+}$ -doped BaGa₂O₄ ceramics, it was shown that the increase of Eu³⁺ content from 1 to 3 mol% leads to agglomeration of free-volume defects near grain boundaries of ceramics. At the same time, nanopores in ceramics expand, and their number increases. Further increase in the content of Eu³⁺ ions is accompanied with fragmentation of both free-volume defects and nanopores.

4. Experimental Section

Samples of the pure BaGa₂O₄ ceramics with Eu³⁺ impurities were prepared by high-temperature solid-phase synthesis similar to spinel-type $MgGa_2O_4$ ceramics.^[49] Powders of BaCO₃ and Ga_2O_3 oxides with a purity of 99.99% were taken as starting components for synthesis. Isopropyl CH₃-CH(OH)-CH₃ alcohol was also used at the sintering of $BaGa_2O_4$ -based ceramics.^[48] The oxide mixtures were taken in a stoichiometric ratio of 1:1 mol. per 2 g of raw material. The content of Eu^{3+} ions (Eu_2O_3 with a purity of 99.99%) was determined to replace the equivalent molar content of Ga and Ba, respectively. The amount of Eu^{3+} ions was 0, 1, 3, and 4 mol%). The oxide mixture was thoroughly stirred in the agate solution for 6 h with the addition of isopropyl alcohol. The obtained raw material was dried in air for 1 h at a temperature of 80 °C. After that, the raw material was pressed under a pressure of $150\,kg\,cm^{-2},$ obtaining workpieces with a diameter of 6 mm and a thickness of 1.5 mm. The blanks were placed on a platinum substrate and annealed in a furnace at 1200 °C for 12 h in air. In the next stat, annealing was performed at 1300 $^\circ\!C$ for 5 h. The final ceramic samples were obtained in the form of tablets with a diameter of 4 mm and a thickness of 1 mm. $\ensuremath{^{[48]}}$

The ORTEC system with ²²Na isotope as positron source was used for PAL measurements. Investigation was performed at 22 $^{\circ}$ C and relative humidity of 35% for two identical ceramic samples placed in a sandwich configuration as described in detail in ref. [42].

The measured PAL spectra were calculated using LT software^[50] involving a three-component fitting procedure (lifetimes τ_1 , τ_2 , τ_3 and intensities l_1 , l_2 , l_3), which is considered a fair approach for spinel ceramics with branched porous structure.^[47,51] The resulting inaccuracy in the positron lifetime of the first component τ_1 was ± 0.002 ns and errors for positron lifetimes of the second and third components (τ_2 and τ_3) were ± 0.001 ns. The error bar for intensities of all component is $\pm 0.1\%$.

Parameters such as the average lifetime of positrons τ_{av} , the lifetime of positrons in defect-free bulk τ_b and the rate of positron trapping in defects κ_d were obtained using a two-stage positron trapping model.^[47] The $\tau_2 - \tau_b$ difference (which describes the size of the free-volume defects where positrons are trapped) and the τ_2/τ_b ratio (correlates with the nature of defects) were also determined. Using the lifetime of the third component, the radii of the nanopores (R_3) were determined according to the Tao–Eldrup model.^[52]

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Conflict of Interest

The authors declare no conflict of interest.

Data Availability Statement

The data that support the findings of this study are available from the corresponding author upon reasonable request.

Keywords

agglomeration, ceramics, fragmentation, free-volume defects, positron annihilation $% \left({{{\left[{{{c_{\rm{s}}}} \right]}_{\rm{s}}}_{\rm{s}}} \right)$

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- L. L. Noto, D. Poelman, V. R. Orante-Barrón, H. C. Swart, L. E. Mathevula, R. Nyenge, M. Chithambo, B. M. Mothudi, M. S. Dhlamini, *Physica B* 2018, 535, 268.
- [2] M. Kodu, T. Avarmaa, H. Mändar, R. Jaaniso, Appl. Phys. A 2008, 93, 801.
- [3] D. Gust, T. A. Moore, A. L. Moore, Acc. Chem. Res. 2009, 42, 1890.
- [4] N. Skillen, P. K. Robertson, in *Solar Energy* (Ed: G. M. Crawley), World Scientific, Singapore **2016**, chapter 6, pp. 205–241.
- [5] W. Acuna, J. F. Tellez, M. A. Macias, P. Roussel, S. Ricote, G. H. Gauthier, Solid State Sci. 2017, 71, 61.
- [6] T. Guérineau, C. Strutynski, T. Skopak, S. Morency, A. Hanafi, F. Calzavara, Y. Ledemi, S. Danto, T. Cardinal, Y. Messaddeq, E. Fargin, Opt. Mater. Express 2019, 9, 2437.
- [7] J. M. Gonçalves, R. A. Munoz, L. Angnes, J. Mater. Chem. 2021, 100, 8842.
- [8] D. Nakauchi, G. Okada, N. Kawaguchi, T. Yanagida, *Opt. Mater.* 2019, 87, 58.
- [9] T. A. Khattab, M. Abd El-Aziz, M. S. Abdelrahman, M. El-Zawahry, S. Kamel, *Luminescence* 2020, 35, 478.
- [10] L. Yu, D. Den Engelsen, K. Gorobez, G. R. Fern, T. G. Ireland, C. Frampton, J. Silver, Opt. Mater. Express 2019, 9, 2175.
- [11] V. M. Maphiri, M. R. Mhlongo, T. T. Hlatshwayo, T. E. Motaung, L. F. Koao, S. V. Motloung, *Opt. Mater.* 2020, 109, 110244.
- [12] D. den Engelsen, G. R. Fern, T. G. Ireland, F. Yang, J. Silver, Opt. Mater. Express 2020, 10, 1962.
- [13] D. den Engelsen, G. R. Fern, T. G. Ireland, J. Silver, ECS J. Solid State Sci. Technol. 2020, 9, 026001.
- [14] N. J. Shivaramu, E. Coetsee, W. D. Roos, K. R. Nagabhushana, H. C. Swart, J. Phys. D: Appl. Phys. 2020, 53, 475305.
- [15] Q. Xie, B. Li, X. He, M. Zhang, Y. Chen, Q. Zeng, Materials 2017, 10, 1198.
- [16] S. Wang, Y. Wang, H. Gao, J. Li, L. Fang, X. Yu, S. Tang, X. Zhao, G. Sun, *Optik* **2020**, *221*, 165363.
- [17] M. Malkamäki, A. J. Bos, P. Dorenbos, M. Lastusaari, L. C. Rodrigues, H. C. Swart, J. Hölsä, *Physica B* **2020**, *593*, 411947.
- [18] D. Nakauchi, G. Okada, T. Kato, N. Kawaguchi, T. Yanagida, *Radiat. Meas.* 2020, 135, 106365.
- [19] M. Vrankić, A. Šarić, S. Bosnar, D. Pajić, J. Dragović, A. Altomare, A. Falcicchio, J. Popović, M. Jurić, M. Petravić, I. Jelovica Badovinac, G. Dražić, *Sci. Rep.* **2019**, *9*, 15158.
- [20] M. G. Brik, A. Suchocki, A. Kaminska, Inorg. Chem. 2014, 53, 5088.
- [21] Y. Wang, W. B. Chen, F. Y. Liu, D. W. Yang, Y. Tian, C. G. Ma, M. D. Dramićanin, M. G. Brik, *Results Phys.* 2019, 13, 102180.
- [22] A. Luchechko, Y. Zhydachevskyy, S. Ubizskii, O. Kravets, A. I. Popov, U. Rogulis, E. Elsts, E. Bulur, A. Suchocki, *Sci. Rep.* 2019, *9*, 9544.
- [23] M. G. Brik, N. M. Avram, C. N. Avram, C. Rudowicz, Y. Y. Yeung, P. Gnutek, J. Alloys Compd. 2007, 432, 61.
- [24] I. Karbovnyk, J. Collins, I. Bolesta, A. Stelmashchuk, A. Kolkevych, S. Velupillai, S. H. Klym, O. Fedyshyn, S. Tymoshuk, I. Kolych, *Nanoscale Res. Lett.* 2015, 10, 151.
- [25] S. V. Lisovskii, A. V. Meganov, V. R. Khrustov, V. V. Ivanov, J. Phys. Conf. Ser. 2019, 1410, 012089.
- [26] D. Valiev, O. Khasanov, E. Dvilis, S. Stepanov, V. Paygin, A. Ilela, Phys. Status Solidi B 2020, 257, 1900471.

- [27] Q. Chen, L. Shang, H. Xu, C. Ma, C. K. Duan, J. Phys. Chem. C 2021, 125, 21780.
- [28] A. Lushchik, E. Feldbach, E. A. Kotomin, I. Kudryavtseva, V. N. Kuzovkov, A. I. Popov, V. Seeman, E. Shablonin, *Sci. Rep.* 2020, *10*, 7810.
- [29] N. Mironova-Ulmane, A. I. Popov, G. Krieke, A. Antuzevics, V. Skvortsova, E. Elsts, A. Sarakovskis, *Low Temp. Phys.* 2020, 46, 1154.
- [30] A. Platonenko, D. Gryaznov, E. A. Kotomin, A. Lushchik, V. Seeman, A. I. Popov, Nucl. Instrum. Methods Phys. Res., Sect. B 2020, 464, 60.
- [31] Q. Li, T. Liu, X. Xu, X. Wang, R. Guo, X. Jiao, Y. Lu, Mater. Technol. 2021, 36, 279.
- [32] Q. Li, T. Liu, X. Xu, R. Guo, X. Jiao, X. Wang, Y. Lu, J. Phys. Chem. Solids 2020, 145, 109542.
- [33] V. Seeman, E. Feldbach, T. Kärner, A. Maaroos, N. Mironova-Ulmane, A. I. Popov, E. Shablonin, E. Vasil'chenko, A. Lushchik, *Opt. Mater.* 2019, *91*, 42.
- [34] O. V. Ogorodnikova, M. Majerle, J. Čížek, S. Simakov, V. V. Gann, P. Hruška, J. Kameník, J. Pospíšil, M. Štefánik, M. Vinš, *Sci. Rep.* 2020, 10, 18898.
- [35] R. Rementeria, R. Domínguez-Reyes, C. Capdevila, C. Garcia-Mateo, F. G. Caballero, *Sci. Rep.* 2020, *10*, 486.
- [36] D. V. Barad, P. L. Mange, K. K. Jani, S. Mukherjee, M. Ahmed, S. Kumar, S. N. Dolia, R. Pandit, P. Y. Raval, K. B. Modi, P. M. G. Nambissan, *Ceram. Int.* **2021**, *47*, 2631.
- [37] H. Klym, A. Ingram, I. Hadzaman, I. Karbovnyk, I. Vasylchyshyn,
 A. I. Popov, IOP Conf. Ser.: Mater. Sci. Eng. 2019, 503, 012019.
- [38] A. H. Ghanem, K. R. Mohamed, Arab J. Nucl. Sci. Appl. 2021, 54, 98.
 [39] H. I. Zhang, S. Sellaiyan, K. Sako, A. Uedono, Y. Taniguchi,
- [39] H. I. Zhang, S. Sellaiyan, K. Sako, A. Uedono, Y. Taniguchi,
 K. Hayashi, *Polymer* **2020**, *190*, 122225.
- [40] D. Biswas, A. Rajan, S. Kabi, A. S. Das, L. S. Singh, P. M. G. Nambissan, *Mater. Charact.* 2019, 158, 109928.
- [41] O. Melikhova, J. Kuriplach, I. Prochazka, J. Cizek, M. Hou, E. Zhurkin, S. Pisov, Appl. Surf. Sci. 2008, 255, 157.
- [42] J. Filipecki, A. Ingram, H. Klym, O. Shpotyuk, M. Vakiv, J. Phys. Conf. Ser. 2007, 79, 012015.
- [43] H. Klym, I. Hadzaman, V. Gryga, Appl. Nanosci. 2021, 12, 1257.
- [44] O. Shpotyuk, V. Balitska, M. Brunner, I. Hadzaman, H. Klym, Physica B 2015, 459, 116.
- [45] H. Klym, A. Ingram, O. Shpotyuk, I. Karbovnyk, Opt. Mater. 2016, 59, 39.
- [46] H. Klym, A. Ingram, O. Shpotyuk, Materialwiss. Werkstofftech. 2016, 47, 198.
- [47] H. Klym, A. Ingram, J. Phys. Conf. Ser. 2007, 79, 012014.
- [48] Y. Kostiv, A. Luchechko, H. Klym, I. Karbovnyk, B. Sadovyi, O. Zaremba, O. Kravets, in XIth Inter. Scientific and Practical Conf. on Electronics and Information Technologies, IEEE, Lviv, Ukraine 2019, p. 307.
- [49] A. Luchechko, O. Kravets, Phys. Status Solidi C 2017, 14, 1600146.
- [50] J. Kansy, D. Giebel, J. Phys. Conf. Ser. 2011, 265, 012030.
- [51] A. Bondarchuk, O. Shpotyuk, A. Glot, H. Klym, *Rev. Mex. Fis.* 2012, 58, 313.
- [52] T. Goworek, Chem. Phys. Lett. 2002, 366, 184.